

**Problem Set 2**  
 $\mathcal{N} = 4$  SYM, BPS property

**Reading:** D'Hoker-Freedman, §2-4; Polchinski vol II appendix B.

**Due:** Thursday, October 2, 2007, roughly.

1. **How to remember the  $\mathcal{N} = 4$  action.**

Show that ten-dimensional  $\mathcal{N} = 1$  SYM dimensionally reduces to 4d  $\mathcal{N} = 4$ . The 10d lagrangian density is

$$\mathcal{L}_{10} = -\frac{1}{2g_{YM}^2} \text{tr} (F_{MN}F^{MN} - 2i\bar{\lambda}\Gamma^M D_M \lambda);$$

$M, N$  are 10d indices;  $\lambda$  here is a 10d Majorana-Weyl spinor (16 real components). By ‘dimensionally reduce,’ we mean consider the 10d theory on a 6-torus of volume  $V$ , and restrict to field configurations which have no momentum along the torus (*i.e.* are independent of the coordinates on the torus).

The 4d  $\mathcal{N} = 4$  lagrangian density is<sup>1</sup>

$$\mathcal{L}_4 = -\frac{1}{g_{YM}^2} \text{tr} \left( \frac{1}{4} F^2 + \frac{1}{2} D_\mu X^i D^\mu X^i + \frac{i}{2} \bar{\lambda}^I \gamma^\mu D_\mu \lambda^I - \frac{1}{2} \sum_{i < j} [X^i, X^j]^2 + \frac{1}{2} \bar{\lambda}^I \Gamma_{IJ}^i [X^i, \lambda^J] \right).$$

Here  $\lambda^I$  are four 4d Weyl spinors (4 real components),  $\gamma^\mu$  are 4d gamma matrices, and  $\Gamma_{IJ}^i$  are gamma matrices for  $SO(6)$ .

2. Show that the  $\mathcal{N} = 4$  lagrangian explicitly breaks the diagonal abelian part

$$U(1) \subset U(\mathcal{N} = 4)_R$$

of the R-symmetry group.

3. Show that the 1-loop  $\beta$ -function for  $\mathcal{N} = 4$  SYM vanishes. If you don't like arcane group theory, do it for the case when the gauge group is  $SU(N)$ .

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<sup>1</sup>This footnote corrects the expression I wrote in lecture (the Yukawa terms were wrong).

4.  $\mathcal{N} = 4 \supset \mathcal{N} = 1$ .

A 4d  $\mathcal{N} = 4$  supersymmetric theory is a special case of a 4d  $\mathcal{N} = 1$  supersymmetric theory.

a) Describe the field content of  $\mathcal{N} = 4$  SYM in terms of multiplets of some  $\mathcal{N} = 1$  subgroup.

b) Write the  $\mathcal{N} = 4$  lagrangian in  $\mathcal{N} = 1$  superspace.

c) Convince yourself that this Lagrangian has  $SO(6)$  R-symmetry (*e.g.* by examining the Lagrangian written in components). Is  $\mathcal{N} = 1$  supersymmetry plus this R-symmetry enough to convince you that this action is actually  $\mathcal{N} = 4$  supersymmetric?

5. **Extremal = BPS.**

In this problem we will consider 4d  $\mathcal{N} = 2$  supergravity (which has eight real supercharges). Along the lines of the discussion from lecture 5 one can see that the graviton multiplet for this theory must also contain a spin-3/2 gravitino field  $\psi_\mu$ , and an abelian vector field (the ‘graviphoton’)  $A_\mu$ . So this is a simple supersymmetric completion of the system we studied in problem set 1 problem 3.

The action is invariant under supersymmetry transformations which include the following variation for the gravitino  $\psi_\mu$ :

$$\delta\psi_\mu = \left( \nabla_\mu - \frac{1}{4}F^- \gamma_\mu \right) \epsilon$$

where

$$F^- \equiv F_{\nu\rho}^- \gamma^\nu \gamma^\rho, \quad F_{\nu\rho}^\pm = \frac{1}{2} \left( F_{\nu\rho} \pm (\star F)_{\nu\rho} \right)$$

are the self-dual (SD) and anti-self-dual (ASD) parts of the field strength of the graviphoton field. The covariant derivative acting on spinors is

$$\nabla_\mu \psi \equiv \left( \partial_\mu + \frac{1}{4} \omega_\mu^{ab} \gamma_{ab} \right) \psi$$

where  $\omega$  is the spin connection,  $de^a = -\omega^{ab} \wedge e^b$  and  $\gamma_{ab} \equiv \frac{1}{2}[\gamma_a, \gamma_b]$ .

a) Show that the extremal RN black hole from problem set 1 lies in a short (BPS) multiplet of the  $\mathcal{N} = 2$  supersymmetry. Do this by examining the supersymmetry variations of the fermionic fields (given above) and showing

that they vanish for some choice of spinors  $\epsilon$ . How many supersymmetries (choices of  $\epsilon$ , which are called Killing spinors) are left unbroken?

b) Show that the ‘Bertotti-Robinson’ solution, *i.e.*  $AdS_2 \times S^2$  (which arises as the near-horizon limit of the extremal RN black hole) preserves eight Killing spinors.

c) If you are feeling energetic, show that  $AdS_5 \times S^5$  preserves 32 supercharges of 10-dimensional type IIB supergravity. The supersymmetry variations of the fermion fields in that case are

$$\delta\psi_\mu = \left( \nabla_\mu - \frac{i}{1920} F_5^- \Gamma_\mu \right) \epsilon + \dots$$

$$\delta\lambda = \dots$$

where

$$F_5^+ \equiv \Gamma^{\mu_1 \dots \mu_5} (F_5)_{\mu_1 \dots \mu_5} \ ,$$

$\lambda$  is the ‘dilatino’ field and ... means terms that vanish when  $F_5$  is the only nontrivial field other than the metric (*i.e.* they depend on derivatives of the dilaton, and the other fluxes). If you want to know what the full variations are, see Kiritsis eqns (H.26-28).

6. **W-bosons from adjoint higgsing.** Using the  $\mathcal{N} = 4$  lagrangian, show that giving an expectation value to the scalars in the  $\mathcal{N} = 4$  theory of the form

$$\langle X^{i=1, \dots, 6} \rangle = \text{diag}(x_1^i, \dots, x_N^i)$$

gives a mass to the off-diagonal gauge bosons of the form

$$m_{ab}^2 \propto \sum_{i=1}^6 (x_a^i - x_b^i)^2 \equiv |\vec{x}_a - \vec{x}_b|^2 \ ;$$

this is meant to be the mass of the gauge boson with indices  $ab$ . Show that the scalars with the same gauge charges get the same mass.