University of California at San Diego – Department of Physics – Prof. John McGreevy

Quantum Mechanics C (Physics 130C) Winter 2015 Assignment 5

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Due 11am Thursday, February 12, 2015

1. More on von Neumann entropy

Consider the general density matrix for a qbit,

$$\boldsymbol{\rho}_{\vec{n}} = \frac{1}{2} \left(\mathbbm{1} + \vec{n} \cdot \boldsymbol{\sigma} \right).$$

Recall that $\vec{n}, \vec{n} \cdot \vec{n} \leq 1$, specifies a point on the *Bloch ball*.

Compute the von Neumann entropy $S(\boldsymbol{\rho}_{\vec{n}})$, and express it as a function of a single variable. Interpret this variable in terms of the geometry of the space of states.

[Hint: use the fact that the von Neumann entropy is independent of the choice of basis.]

2. Entropy and thermodynamics

Consider a quantum system with hamiltonian **H** and Hilbert space $\mathcal{H} = \text{span}\{|n\rangle, n = 0, 1, 2...\}$. Its behavior at finite temperature can be described using the *thermal density* matrix

$$\boldsymbol{
ho}_{eta} \equiv rac{1}{Z} e^{-eta \mathbf{H}}$$

where $\beta \equiv \frac{1}{T}$ specifies the temperature and Z is a normalization factor. (We can think about this as the reduced density matrix resulting from coupling the system to a heat bath and tracing out the Hilbert space of the (inaccessible) heat bath.)

- (a) Find a formal expression for Z by demanding that ρ_{β} is normalized appropriately. (Z is called the *partition function*.)
- (b) Show that the von Neuman entropy of ρ_{β} can be written as

$$S_{\beta} = E/T + \log Z \tag{1}$$

where $E \equiv \langle \mathbf{H} \rangle_{\boldsymbol{\rho}} = \text{tr} \boldsymbol{\rho} \mathbf{H}$ is the expectation value for the energy. The expression above for S_{β} is the thermal entropy.

(c) Evaluate Z and E for the case where the system is a simple harmonic oscillator

$$\mathbf{H} = \hbar\omega \left(\mathbf{n} + \frac{1}{2}\right)$$

with $\mathbf{n}|n\rangle = n|n\rangle$.

[For further inquiry in this direction I recommend Schumacher and Westmoreland sections 19.5, 19.6.]

3. Purity test. [from Chuang and Nielsen]

Show that for any density matrix ρ :

- (a) $tr \rho^2 \leq 1$
- (b) the inequality is saturated only if ρ is a pure state.

[Hint: don't forget that the trace operation is basis-independent.]

4. Polarization. [from Boccio]

Recall that in the expression for the general density matrix of a qbit

$$\boldsymbol{\rho} = \frac{1}{2} \left(\mathbb{1} + \vec{P} \cdot \vec{\boldsymbol{\sigma}} \right)$$

we called \vec{P} the *polarization*.

(a) To justify this name, show that

$$\vec{P} = \langle \vec{\sigma} \rangle.$$

(b) Subject the qbit to a external magnetic field, which couples by

$$\mathbf{H} = -\vec{\boldsymbol{\mu}} \cdot \vec{B} = -\frac{\gamma}{2} \boldsymbol{\sigma} \cdot \vec{B}.$$

(γ is called the *gyromagnetic ratio*.) Assuming the qbit is isolated, so that its time evolution is unitary, what is the time evolution of the polarization, $\partial_t \vec{P}$?

5. Quantum gates

Schumacher and Westmoreland, Exercises 18.2, 18.6, 18.7, and 18.8, page 368. (Electronic version here.)

Regarding the SWAP gate, Schumacher's wording is confusing and basis-dependent. This gate is defined to act on two qbits by

$$\mathbf{S} = |00\rangle\langle00| + |01\rangle\langle10| + |10\rangle\langle01| + |11\rangle\langle11|$$

(so in words, it acts by swapping the labels of the basis states of the tensor-product Hilbert space). This implies (check this) that acting on *product* (*i.e.* unentangled) states $|a\rangle \otimes |b\rangle$ it gives $\mathbf{S}|a\rangle \otimes |b\rangle = |b\rangle \otimes |a\rangle$.