University of California at San Diego – Department of Physics – Prof. John McGreevy

Physics 215A QFT Fall 2017 Assignment 5

Due 11am Thursday, November 2, 2017

1. Recovering non-relativistic quantum mechanics.

Consider a complex scalar field, in the non-relativistic limit,

$$\Phi = \sqrt{2m}e^{-\mathbf{i}mt}\Psi, \quad |\dot{\Psi}| \ll m\Psi.$$

Recall that in this limit, the antiparticles disappear and the mode expansion is

$$\Psi(x) = \int \mathrm{d}^d p \ e^{+\mathbf{i}\vec{p}\cdot\vec{x}} \mathbf{a}_p, \quad \Psi^{\dagger}(x) = \int \mathrm{d}^d p \ e^{-\mathbf{i}\vec{p}\cdot\vec{x}} \mathbf{a}_p^{\dagger}$$

(a) Show that

$$\hat{P}_i \equiv \int \mathrm{d}^d p p_i \mathbf{a}_p^\dagger \mathbf{a}_p$$

is the generator of translations and commutes with the Hamiltonian.

(b) Let

$$\hat{X}^i \equiv \int d^d x \Psi^{\dagger}(x) x^i \Psi(x).$$

A state of one particle at location \vec{x} is

$$|x\rangle = \Psi^{\dagger}(x) |0\rangle.$$

Show that

$$\hat{X}^{i} \left| x \right\rangle = x^{i} \left| x \right\rangle.$$

(c) Consider the general one-particle state

$$|\psi\rangle = \int d^d x \ \psi(x) \Psi^{\dagger}(x) \left|0\right\rangle = \int d^x x \ \psi(x) \left|x\right\rangle.$$

Show that

$$\hat{X}^{i} \left| \psi \right\rangle = \int d^{d}x x^{i} \psi(x) \left| x \right\rangle$$

and (a little more involved)

$$\hat{P}^{i} \left| \psi \right\rangle = \int d^{d}x \left(-\mathbf{i} \frac{\partial}{\partial x^{i}} \psi(x) \right) \left| x \right\rangle,$$

which is the usual action of these operators on single-particle wavefunctions $\psi(x)$.

2. Scalar Yukawa amplitudes.

Consider again the scalar Yukawa theory of a complex scalar Φ and a real scalar ϕ . In the following, assume all particles are in momentum eigenstates. Use artisanal methods.

- (a) Compute the amplitude for the annihilation of a Φ particle and a Φ^* particle into a ϕ particle, at leading order in the coupling g.
- (b) Compute the amplitude for $\Phi + \phi \rightarrow \Phi + \phi$ scattering to the leading order in the coupling at which it is nonzero.

3. Wick example.

For a real scalar field, verify by hand Wick's prediction for the difference

$$\mathcal{T}(\phi(x_1)\phi(x_2)\phi(x_3)) - : \phi(x_1)\phi(x_2)\phi(x_3) :$$

- 4. Fields and forces. [from Banks] Consider a real free relativistic scalar field of mass $m S[\phi] = \int d^{d+1}x \frac{1}{2} (\partial_{\mu}\phi \partial^{\mu}\phi m^2\phi^2).$
 - (a) Calculate the vacuum expectation value

$$\langle 0 | \mathcal{T} \left(e^{\mathbf{i} \int d^{d+1}x \ \phi(x)J(x)} \right) | 0 \rangle \equiv e^{\mathbf{i}W[J]}$$

where J is a fixed, external source. Use Wick's theorem. Make a series expansion in powers of J and draw some diagrams. To understand the structure of the series, recall the formula on a previous homework for $\langle e^{K \cdot q} \rangle$ in any gaussian theory.

(b) Now specialize to the case where the source is static and is present for a time 2T:

$$J(x) = J_{\text{static}} \equiv \theta(T-t)\theta(t+T) \left(\delta^{d}(x) - \delta^{d}(x-R)\right)$$

with $T \gg R \gg 1/m$. You should find an answer of the form

$$W[J_{\text{static}}(x)] = TV(R)$$

where V(R) is the Yukawa potential.

(c) Chant the following incantation:

Static sources experience a force due to exchange of virtual particles.

Feel happy at having reproduced by canonical methods the answer we found earlier using path integral methods.