

Physics 212C QM Spring 2020 Assignment 4

Due 12:30pm Monday, April 27, 2020

Please refer to the first homework for submission format and procedures (and replace hw01 by hw04 in the relevant places, of course).

1. **Commutation relations of creation operators for general one-particle states.** Show that

$$\mathbf{a}(\varphi_1)\mathbf{a}^\dagger(\varphi_2) - \zeta\mathbf{a}^\dagger(\varphi_2)\mathbf{a}(\varphi_1) = \langle\varphi_2|\varphi_1\rangle,$$

where these objects are as defined in the lecture notes.

2. **Fermion creation and annihilation algebra.**

Consider a single fermion mode \mathbf{c} . We showed in lecture that the associated hilbert space is two-dimensional, and is spanned by

$$|0\rangle, \text{ with } \mathbf{c}|0\rangle = 0 \text{ and } |1\rangle = \mathbf{c}^\dagger|0\rangle.$$

- (a) Check that the two states are orthogonal.
- (b) Show that acting on this Hilbert space it is indeed true that

$$\mathbf{c}^\dagger\mathbf{c} + \mathbf{c}\mathbf{c}^\dagger = \mathbb{1},$$

as long as $\langle 1|1\rangle = \langle 0|0\rangle$.

- (c) Check that

$$[\mathbf{N}, \mathbf{c}] = -\mathbf{c}, \quad [\mathbf{N}, \mathbf{c}^\dagger] = \mathbf{c}^\dagger$$

where $\mathbf{N} = \mathbf{c}^\dagger\mathbf{c}$ is the number operator. Notice that this is the same algebra satisfied by bosonic modes.

3. **Majorana modes.** Given a collection of fermionic operators \mathbf{c}_A , satisfying the fermionic creation-annihilation algebra

$$\{\mathbf{c}_A, \mathbf{c}_B^\dagger\} = \delta_{AB}\mathbb{1} \quad \text{and} \quad \{\mathbf{c}_A, \mathbf{c}_B\} = 0,$$

we can decompose them into their real and imaginary parts

$$\gamma_{A1} \equiv \frac{1}{2}(\mathbf{c}_A + \mathbf{c}_A^\dagger), \quad \gamma_{A2} \equiv \frac{1}{2i}(\mathbf{c}_A - \mathbf{c}_A^\dagger).$$

These are called *Majorana modes*.

(a) Show that the Majorana modes satisfy the algebra

$$\{\gamma_a, \gamma_b\} = 2\Upsilon\delta_{ab}\mathbb{1},$$

where here a is a multi-index running over both A and $\alpha = 1, 2$. In particular, notice that $\gamma_a^2 = \Upsilon\mathbb{1}$. Find the constant Υ .

(b) Write the number operator $\mathbf{c}_A^\dagger \mathbf{c}_A$ in terms of the Majorana modes. Show that it is hermitian.

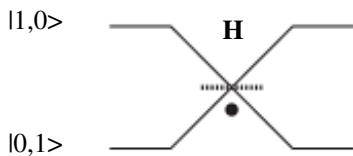
4. **Normalization.** Check that if $\Psi(r_1 \cdots r_n)$ is a normalized and (anti)symmetric wavefunction on n particles, then

$$|\Psi\rangle \equiv \sum_{r_1 \cdots r_n} \Psi(r_1 \cdots r_n) |r_1 \cdots r_n\rangle \quad (1)$$

is normalized, $\langle\Psi|\Psi\rangle = 1$.

(Interpret the sum over r as an integral if you like.)

5. **Multiple photons on paths of an interferometer.**



One way to make a qbit is out of the two states of a photon moving on the upper and lower paths of an interferometer. On such a qbit, a half-silvered mirror \mathbf{H} acts as a unitary gate, as indicated at left. (The dot below the mirror specifies a sign convention, to be explained below.)

On the other hand, photons are bosons. This means that if

$$\mathbf{a}^\dagger |0, 0\rangle \equiv |1, 0\rangle \text{ is a state with one photon on the upper path}$$

of the interferometer, then

$$\frac{(\mathbf{a}^\dagger)^n}{\sqrt{n!}} |0, 0\rangle \equiv |n, 0\rangle \text{ is a state with } n \text{ photons on the upper path.}$$

Similarly, define

$$\frac{(\mathbf{b}^\dagger)^n}{\sqrt{n!}} |0, 0\rangle \equiv |0, n\rangle \text{ to be a state with } n \text{ photons on the lower path}$$

of the interferometer. (Note that $[\mathbf{a}, \mathbf{b}] = 0 = [\mathbf{a}, \mathbf{b}^\dagger]$ – they are independent modes.)

Now suppose we direct these two paths through a half-silvered mirror, as in the figure. A half-silvered mirror acts as a Hadamard gate

$$\mathbf{H} \equiv \frac{1}{\sqrt{2}} (\boldsymbol{\sigma}^x + \boldsymbol{\sigma}^z) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

on the qbit made from the one-photon states. (The dot tells us where to put the negative entry.)

Some warm-up questions:

- (a) What is the state $|0, 0\rangle$? How does \mathbf{H} act on $|0, 0\rangle$?
- (b) How does \mathbf{H} act on $|2, 0\rangle$ and $|0, 2\rangle$?
- (c) How does \mathbf{H} act on the operators \mathbf{a}^\dagger and \mathbf{b}^\dagger (in order that the above relations are realized)?

Here's a more interesting question:

- (d) What is the state which results upon sending a coherent state of photons

$$|\alpha, \beta\rangle \equiv \mathcal{N}_\alpha \mathcal{N}_\beta e^{\alpha \mathbf{a}^\dagger + \beta \mathbf{b}^\dagger} |0, 0\rangle$$

through a half-silvered mirror? ($\mathcal{N}_\alpha \equiv e^{-|\alpha|^2/2}$ is a normalization constant.)

[Hint: it may be useful to insert $\mathbb{1} = \mathbf{H}^2$ in between the $e^{\alpha \mathbf{a}^\dagger + \beta \mathbf{b}^\dagger}$ and the $|0, 0\rangle$.]

The half-silvered mirror is a special case of the more general notion called a beam-splitter. Suppose instead that the action on the mode operators were

$$\mathbf{U}^\dagger(\theta) \mathbf{a} \mathbf{U}(\theta) = \mathbf{a} \cos \theta + \mathbf{i} \mathbf{b} \sin \theta \quad (2)$$

$$\mathbf{U}^\dagger(\theta) \mathbf{b} \mathbf{U}(\theta) = \mathbf{b} \cos \theta + \mathbf{i} \mathbf{a} \sin \theta . \quad (3)$$

- (e) Show that $\mathbf{U}(\theta)$ can be written as an evolution operator, in the form:

$$\mathbf{U}(\theta) = e^{i\theta G}, \quad G = \mathbf{a}^\dagger \mathbf{b} + \mathbf{b}^\dagger \mathbf{a}.$$

- (f) Show that when $\theta = \pi/4$ this beam-splitter takes the state $|1, 1\rangle$ with one boson in each mode to the state

$$\frac{1}{\sqrt{2}} (|2, 0\rangle + |0, 2\rangle).$$

- (g) What if the operators \mathbf{a} and \mathbf{b} were instead fermionic operators? That is, suppose we send fermionic particles through the same beam-splitter. What is

$$\mathbf{U}_F(\theta = \pi/4)^\dagger |1, 1\rangle$$

in this case? Hint: the form of the generator is different

$$\mathbf{U}_F(\theta) = e^{i\theta G_F}, \quad G_F = \mathbf{a}^\dagger \mathbf{b} - \mathbf{a} \mathbf{b}^\dagger.$$

(Notice that G_F is still hermitian.)