University of California at San Diego - Department of Physics - Prof. John McGreevy

## Physics 215C QFT Spring 2022 Assignment 5

Due 11:59pm Monday, May 2, 2022
Thanks in advance for following the guidelines on HW01. Please ask me by email if you have any trouble.

## 1. Boson coherent states brain warmers.

Verify the following identities for the coherent state $|\phi\rangle=e^{\phi \mathbf{a}^{\dagger}}|0\rangle$ of a single mode.
(a)

$$
\left\langle\phi_{1} \mid \phi_{2}\right\rangle=e^{\phi_{1}^{\star} \phi_{2}} .
$$

(b)

$$
\mathbb{1} \equiv \sum_{n=0}^{\infty}|n\rangle\langle n|=\int \frac{d \phi d \phi^{\star}}{\pi} e^{-|\phi|^{2}}|\phi\rangle\langle\phi| .
$$

(c)

$$
\operatorname{tr} \cdot=\int \frac{d \phi d \phi^{\star}}{\pi} e^{-|\phi|^{2}}\langle\phi| \cdot|\phi\rangle .
$$

## 2. Grassmann exercises.

(a) A useful device is the integral representation of the grassmann delta function. Show that

$$
-\int d \bar{\psi}_{1} e^{-\bar{\psi}_{1}\left(\psi_{1}-\psi_{2}\right)}=\delta\left(\psi_{1}-\psi_{2}\right)
$$

in the sense that $\int d \psi_{1} \delta\left(\psi_{1}-\psi_{2}\right) f\left(\psi_{1}\right)=f\left(\psi_{2}\right)$ for any grassmann function $f$. (Notice that since the grassmann delta function is not even, it matters on which side of the $\delta$ we put the function: $\int d \psi_{1} f\left(\psi_{1}\right) \delta\left(\psi_{1}-\psi_{2}\right)=f\left(-\psi_{2}\right) \neq$ $f\left(\psi_{2}\right)$.)
(b) Recall the resolution of the identity the Hilbert space of a single fermion mode in terms of fermion coherent states

$$
\begin{equation*}
\mathbb{1} \equiv \sum_{n=0}^{1}|n\rangle\langle n|=\int d \bar{\psi} d \psi e^{-\bar{\psi} \psi}|\psi\rangle\langle\bar{\psi}| . \tag{1}
\end{equation*}
$$

Show that $\mathbb{1}^{2}=\mathbb{1}$. (The previous part may be useful.)
(c) In lecture I claimed that a representation of the trace of a bosonic operator was

$$
\operatorname{tr} \mathbf{A}=\int d \bar{\psi} d \psi e^{-\bar{\psi} \psi}\langle-\bar{\psi}| \mathbf{A}|\psi\rangle
$$

and the minus sign in the bra had important consequences.
(Here $\langle-\bar{\psi}| \mathbf{c}^{\dagger}=\langle-\bar{\psi}|(-\bar{\psi})$ ).
Check that using this expression you get the correct answer for

$$
\operatorname{tr}\left(a+b \mathbf{c}^{\dagger} \mathbf{c}\right)
$$

where $a, b$ are ordinary numbers.
(d) Prove the identity (1) by expanding the coherent states in the number basis.

## 3. Fermionic coherent state exercise.

Consider a collection of fermionic modes $c_{i}$ with quadratic hamiltonian $H=$ $\sum_{i j} h_{i j} c_{i}^{\dagger} c_{j}$, with $h=h^{\dagger}$.
(a) Compute tre $e^{-\beta H}$ by changing basis to the eigenstates of $h_{i j}$ (the singleparticle hamiltonian) and performing the trace in that basis: tr... $=\prod_{\epsilon} \sum_{n_{\epsilon}=c_{\epsilon}^{\dagger} c_{\epsilon}=0,1} \cdots$
(b) Compute tre $e^{-\beta H}$ by coherent state path integral. Compare!
[Hint: to do the Matsubara sum, it is helpful to use an integral representation such as

$$
\sum_{n} f\left(\mathbf{i} \omega_{n}\right)=\frac{1}{2 \pi \mathbf{i}} \oint_{C} \frac{\beta d z}{e^{\beta z}+1} f(z)
$$

where $C$ is a contour that encircles all the poles of $\frac{1}{e^{\beta z}+1}$.]
(c) [super bonus problem] Consider the case where $h_{i j}$ is a random matrix. What can you say about the thermodynamics?

## 4. Topological terms in QM.

The purpose of this problem is to demonstrate that total derivative terms in the action (like the $\theta$ term in QCD) do affect the physics.
The euclidean path integral for a particle on a ring with magnetic flux $\theta=\int \vec{B} \cdot \mathrm{~d} \vec{a}$ through the ring is given by

$$
Z=\int[D \phi] e^{-\int_{0}^{\beta} \mathrm{d} \tau\left(\frac{m}{2} \dot{\phi}^{2}-\mathbf{i} \frac{\theta}{2 \pi} \dot{\phi}\right)}
$$

Here

$$
\begin{equation*}
\phi \equiv \phi+2 \pi \tag{2}
\end{equation*}
$$

is a coordinate on the ring. Because of the identification (6), $\phi$ need not be a single-valued function of $\tau$ - it can wind around the ring. On the other hand, $\dot{\phi}$ is single-valued and periodic and hence has an ordinary Fourier decomposition. This means that we can expand the field as

$$
\begin{equation*}
\phi(\tau)=\frac{2 \pi}{\beta} Q \tau+\sum_{\ell \in \mathbb{Z} \backslash 0} \phi_{\ell} e^{\mathrm{i} \frac{2 \pi}{\beta} \ell \tau} \tag{3}
\end{equation*}
$$

(a) Show that the $\dot{\phi}$ term in the action does not affect the classical equations of motion. In this sense, it is a topological term.
(b) Using the decomposition (7), write the partition function as a sum over topological sectors labelled by the winding number $Q \in \mathbb{Z}$ and calculate it explicitly.
[Hint: use the Poisson resummation formula

$$
\sum_{n} f(n)=\sum_{l} \hat{f}(2 \pi l)
$$

where $\hat{f}(p)=\int d x e^{-\mathbf{i} p x} f(x)$ is the fourier transform of $f$.]
(c) Use the result from the previous part to determine the energy spectrum as a function of $\theta$.
(d) Derive the canonical momentum and Hamiltonian from the action above and verify the spectrum.
(e) Consider what happens in the limit $m \rightarrow 0, \theta \rightarrow \pi$ with $X \equiv \frac{\theta-\pi}{m} \sim \beta^{-1}$ fixed. Interpret the result as the partition function for a spin $1 / 2$ particle. What is the meaning of the ratio $X$ in this interpretation?

