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## Physics 239 Topology from Physics Winter 2021 Assignment 2 – Solutions

Due 12:30pm Monday, January 18, 2020

Thanks in advance for following the guidelines on hw01. Please ask me by email if you have any trouble.

- 1. Toric code as  $\mathbb{Z}_2$  gauge theory with matter. Consider a model with qubits on the links of a lattice (with Pauli operators  $X_{\ell}, Z_{\ell}, X_{\ell}Z_{\ell} = -Z_{\ell}X_{\ell}$ ) and qubits on the sites of the lattice (with Pauli operators  $\sigma_j^x, \sigma_j^z, \sigma_j^x\sigma_j^z = -\sigma_j^z\sigma_j^x$ ).
  - (a) Show that the operator

$$G_j \equiv A_j \sigma_j^z$$

(where  $A_j$  is the star operator) generates the gauge transformation

$$\sigma_j^x \to (-1)^{s_j} \sigma_j^x, \quad X_{ij} \to (-1)^{s_i} X_{ij} (-1)^{s_j}, \quad \sigma_j^z \to \sigma_j^z, \quad Z_{ij} \to Z_{ij}$$
(1)

(where i, j are the sites at the ends of the link labelled ij). By generates here I mean that an operator  $\mathcal{O}$  transforms as

$$\mathcal{O} \to \mathcal{G}_s^{\dagger} \mathcal{O} \mathcal{G}_s, \ \ \mathcal{G}_s \equiv \prod_j G_j^{s_j}$$

with  $s_j = 0, 1$ .

The relations involving Z and  $\sigma^z$  follow because they commute with  $G_j$ .  $X_\ell$  transforms under  $G_j$  if the site  $j \in \partial \ell$ .

(b) Show that the Hamiltonian

$$\mathbf{H} = -\sum_{j} G_{j} - \sum_{p} B_{p} - h \sum_{ij} \sigma_{i}^{x} X_{ij} \sigma_{j}^{x} - g \sum_{\langle ij \rangle} Z_{ij}$$

is gauge invariant.

 $G_j$  commutes with  $G_{j'}$ .  $B_p$  commutes since it is a closed loop of Xs. The third term is a kinetic term for the *e* particles, which is invariant because the transformation of the  $\sigma^x$ s cancels that of the X. The last term is invariant because it is made of Zs.

Here we can identify  $\sigma_j^x$  as the operator which creates an *e* particle at site *j*. And we can identify  $\sigma_j^z = (-1)^{n_j}$  as the parity of the number operator.

(c) Show that if we set  $\sigma_j^x = 1$  and  $\sigma_j^z = 1$  for all j we get back the (perturbed) toric code.

Bonus problem: interpret this operation as a choice of gauge in the model where  $G_j = 1$  is imposed as a constraint on physical states.

Clearly if we just erase all the  $\sigma_j^x$ s and  $\sigma_j^z$ s we get back the toric code Hamiltonian.

And clearly we can choose the gauge parameter  $s_j$  in (1) to set  $\sigma^x = 1$ . The tricky part of this is that we can also erase all the appearances of  $\sigma^z$ . This wouldn't make sense on the full Hilbert space  $\otimes \mathcal{H}_2$ , since  $\sigma^z$  and  $\sigma^x$  do not commute. The claim is that on the space of physical states of the gauge theory, which satisfy  $G_j |\text{phys}\rangle = |\text{phys}\rangle$  for all j, we can do this. It's because on such states, the action of  $\sigma_j^z$  can be replaced by  $A_j$ . If everything is gauge-invariant, any appearance of  $G_j$  can be moved onto the states, and replaced with 1.

2. **3-ball.** Find two cellulations of the 3-dimensional ball (*e.g.* the region  $x^2 + y^2 + z^2 \le 1$  in  $\mathbb{R}^3$ ) and compute the resulting homology groups.

One cellulation is one 3-cell and then the cellulation of the boundary 2-sphere by a single 2-cell and a single 0-cell. The boundary maps are  $\partial \sigma_3 = \sigma_2$ ,  $\partial \sigma_2 = 0$  and we find  $H_i(B_3, \mathbb{Z}) = \delta_{i,0}\mathbb{Z}$ .

Alternatively, we could take a different cellulation of the boundary 2-sphere, say the iterative scheme described in the lecture notes. The cell complex is then

$$0 \to A \xrightarrow{(1,-1)} A^2 \xrightarrow{\begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix}} A^2 \xrightarrow{\begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}} A^2 \to A^2 \to 0.$$

Compared to the cell complex for  $S^2$ , the extra step kills the 2d homology, and we again find  $H_i(B_3, \mathbb{Z}) = \delta_{i,0}\mathbb{Z}$ .

## 3. Who am I?

(a) Compute the homology of the following cell complex: take a 2-simplex  $[v_0, v_1, v_2]$  (recall that the simplex  $[v_0, \dots, v_n]$  is the set of all convex combinations of the points  $v_i$ :  $[v_0, \dots, v_n] \equiv \{\lambda_i v_i, \sum_i \lambda_i = 1, \lambda_i \geq 0\}$ ) and identify the edges  $[v_0, v_1]$  and  $[v_1, v_2]$  (with the orientation preserving the order of the vertices). Also identify the 0-cells in their boundaries. In order to identify the edges  $[v_0, v_1]$  and  $[v_1, v_2]$ , we have to identify all the

vertices. The edge  $[v_2, v_0]$  is still distinct, so there are two 2-cells,  $y_1, y_2$ .

Here is a picture of the resulting cell complex:



The cell complex is

$$0 \to A \xrightarrow{(2,1)} A^2 = \langle y_1 = [v_0, v_1] = [v_1, v_2], [v_2, v_0] \rangle \xrightarrow{\begin{pmatrix} 0 \\ 0 \end{pmatrix}} A \to 0.$$

For  $A = \mathbb{Z}$ , this gives  $H_2 = 0, H_1 = \mathbb{Z} \oplus \mathbb{Z}_2, H_0 = \mathbb{Z}$ .

(b) Compute the homology of the following cell complex: take a square. Identify one pair of opposite edges with a twist:



The other pair of sides remains distinct. Is this the same space from the previous part? What space is it?

Both of these spaces describe a Mobius strip, an unoriented 2d manifold with a single boundary. This cell complex is

$$0 \to \mathbb{Z} \stackrel{(2,1,1)}{\to} \mathbb{Z}^3 \stackrel{\begin{pmatrix} -1 & 1 \\ -1 & 1 \\ 1 & -1 \end{pmatrix}}{\to} \mathbb{Z}^2 \to 0$$

whose homology is

$$H_2 = 0, H_1 = \langle y_1, y_2 + y_3 | 2y_1 = 0 \rangle = \mathbb{Z}_2 \oplus \mathbb{Z}, H_2 = \mathbb{Z} = \langle p_1 = p_2 \rangle.$$

4. Consider a sphere with an extra 1-cell attaching the north pole to the south pole. Compute the homology of this space.

Let's decompose the sphere as a single 2-cell w whose boundary is a single line of longitude,  $e_1$ , traversed twice:  $\partial w = e_1 - e_1 = 0$ . Then the extra 1-cell is  $e_2$ with  $\partial e_1 = N - S = \partial e_2$ . The complex is

$$0 \to A \xrightarrow{0} A^2 \xrightarrow{1 - 1} A^2 \to 0.$$

 $\partial_2$  clearly has rank 1, so we find

$$H_2 = A, H_1 = A = \langle e_1 - e_2 \rangle, H_0 = A = \langle N = S \rangle.$$

Later we'll see that homology is invariant under homotopy. We can find a homotopy that moves the north pole end of the extra 1-cell down to the south pole. So this space has the same homology as a sphere with a loop attached to it.